

ON THE VECTOR FIBER SURFACE OF THE SPACE $Lm(Vn)$ OF TRIPLET CONNECTEDNESS

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Abstract. In this paper, we consider the theory of the surface of metric vector fibers for the space $Lm(Vn)$ with triplet connectedness. It is proved that in metric vector fibers there always exists an internal triplet connectedness. Analogues of Gauss–Weingarten derivation formulas and also analogues of generalized Gauss, Peterson–Codazzi–Mainardi equations are found.

Let us consider the vector fiber space $Lm(Vn)$, where the local coordinates of a point transform by the law [2]

$$\begin{aligned} \bar{x}^i &= \bar{x}^l(x^k); & \bar{y}^\alpha &= A_\beta^\alpha(x)y^\beta; \\ \det \left\| \frac{\partial \bar{x}^i}{\partial x^k} \right\| &\neq 0; & \det \|A_\beta^\alpha\| &\neq 0; & i, j, k &= 1, \dots, n; & \alpha, \beta, \gamma &= 1, \dots, m. \end{aligned} \quad (1)$$

Assume that the tensor field $G_{AB}(A, B, C = 1, 2, \dots, n + m)$ is given on the space $Lm(Vn)$, i.e.

$$\overline{G_{AB}} = \bar{x}_A^C \bar{x}_B^D G_{CD},$$

where

$$\bar{x}_B^A = \left\| \frac{\partial \bar{x}^A}{\partial \bar{x}^B} \right\| = \left\| \begin{array}{cc} \frac{\partial \bar{x}^i}{\partial x^j} & \frac{\partial \bar{x}^i}{\partial x^\alpha} \\ \frac{\partial \bar{x}^\alpha}{\partial x^i} & \frac{\partial \bar{x}^\alpha}{\partial x^\beta} \end{array} \right\| = \left\| \begin{array}{cc} x_j^i & 0 \\ A_{\beta k}^\alpha y^\beta & A_\beta^\alpha \end{array} \right\|.$$

An inverse matrix of the matrix has the form

$$\bar{x}_B^* = \left\| \frac{\partial \bar{x}^A}{\partial \bar{x}^B} \right\| = \left\| \begin{array}{cc} \frac{\partial x^i}{\partial \bar{x}^k} & \frac{\partial x^i}{\partial \bar{x}^\alpha} \\ \frac{\partial x^\alpha}{\partial \bar{x}^i} & \frac{\partial x^\alpha}{\partial \bar{x}^\beta} \end{array} \right\| = \left\| \begin{array}{cc} \bar{x}_j^i & 0 \\ A_{\beta k}^* A_\gamma^\beta y^\gamma & A_\beta^* \end{array} \right\|.$$

Since $G^{\alpha\beta} = A_\gamma^\alpha A_\delta^\beta G^{\gamma\delta}$, where $G_{\beta\gamma} G^{\alpha\beta} = \delta_\gamma^\alpha$ and $G_{\beta i} = A_\beta^\gamma x_i^k G_{\gamma k} + A_\beta^\gamma A_{\varepsilon i}^* A_\rho^\varepsilon y^\rho G_{\gamma \alpha}$, we can use them to construct the values Γ_i^α as follows: $\Gamma_i^\alpha = G^{\alpha\beta} G_{\beta i}$.

Furthermore,

$$G^{\alpha\beta} G_{\beta i} = A_\gamma^\alpha A_\delta^\beta G^{\gamma\delta} A_\beta^* x_i^k G_{\rho k} + A_\gamma^\beta G^{\gamma\delta} A_\beta^\omega A_{\varepsilon i}^* A_\rho^\varepsilon y^\rho G_{\omega \sigma}.$$

Since $A_\gamma^* A_\beta^\alpha = \delta_\beta^\alpha$, $A_\gamma^k A_\varepsilon^\alpha A_\varepsilon^\gamma + A_\gamma^\alpha A_{\varepsilon i}^* x_i^k = 0$, $-x_i^k A_{\gamma k}^* A_{\gamma k}^\alpha A_\varepsilon^\gamma = A_\gamma^\alpha A_{\varepsilon i}^*$, we observe that the values Γ_i^α form an object of linear connectedness with the following transformation law

$$\Gamma_i^\alpha = A_\gamma^\alpha x_i^k \Gamma_k^\gamma - x_i^k A_{\gamma k}^* A_{\gamma k}^\alpha y^\gamma,$$

and the values

$$g_{ij} = G_{ij} - \Gamma_i^\alpha G_{\alpha j} - \Gamma_j^\beta G_{\beta i} + \Gamma_i^\alpha \Gamma_j^\beta G_{\alpha\beta}$$

are a double covariant symmetric tensor so that we can construct an object of affine connectedness Γ_{jk}^i in the following manner

$$\Gamma_{jk}^i - \frac{1}{2}g^{ip}(\nabla_k g_{pj} + \nabla_j g_{kp} - \nabla_p g_{jk}).$$

Note that the linear connectedness Γ_i^α induces the vertical affine connectedness defined by the object $\nabla_\beta \Gamma_i^\alpha \equiv \Gamma_{\beta i}^\alpha$ with the following transformation law

$$\overline{\Gamma_{\beta i}^\alpha} = x_i^k A_\beta^\gamma A_\delta^\alpha \Gamma_{\gamma k}^\delta - A_{\beta i}^\alpha.$$

Structural equations of the space $Lm(Vn)$ with triplet connectedness have the form [3, 4]:

$$\begin{cases} D\omega^i = \omega^k \wedge \tilde{\omega}_k^i, \\ D\tilde{\theta}^\alpha = \tilde{\theta}^\beta \wedge \tilde{\omega}_\beta^\alpha + R_{ik}^\alpha \omega^i \wedge \omega^k, \\ D\tilde{\omega}_\beta^\alpha = \tilde{\omega}_\beta^\gamma \wedge \tilde{\omega}_\gamma^\alpha + R_{\beta ik}^\alpha \omega^i \wedge \omega^k + R_{\beta i \gamma}^\alpha \omega^i \wedge \tilde{\theta}^\gamma, \\ D\tilde{\omega}_j^i = \tilde{\omega}_j^k \wedge \tilde{\omega}_k^i + R_{j pq}^i \omega^p \wedge \omega^q + R_{j p \gamma}^i \omega^p \wedge \tilde{\theta}^\gamma. \end{cases} \quad (2)$$

Assume that a hypersurface \mathfrak{N} is given on the space $Lm(Vn)$

$$\omega^i = M_a^i \psi^a \quad (3)$$

and the 1-forms ψ^a re such that

$$\begin{cases} D\psi^a = \psi^b \wedge \psi_b^a, \\ D\psi_b^a = \psi_b^c \wedge \psi_c^a + \psi^c \wedge \psi_{bc}^a. \end{cases}$$

Note that

$$D\tilde{\theta}^\alpha = \tilde{\theta}^\beta \wedge \tilde{\omega}_\beta^\alpha + R_{ik}^\alpha \omega^i \wedge \omega^k = \tilde{\theta}^\beta \wedge \tilde{\omega}_\beta^\alpha + R_{ik}^\alpha M_a^i \psi^a \wedge M_a^k \psi^a = R_{ab}^\alpha \psi^a \wedge \psi^b,$$

where $R_{ab}^\alpha = R_{ik}^\alpha M_a^i M_b^k$.

The extension of system (3) is given by

$$\begin{cases} \nabla M_a^i = M_{ab}^i \psi^b, & \nabla M_{ab}^i + M_c^i \psi_{ab}^c = M_{abc}^i \psi^c, \\ \nabla M_{abc}^i + 2M_{(a|d|}^i \psi_{b)c}^d - M_d^i \psi_{bc}^d = M_{abd}^i \psi^d, \end{cases}$$

where

$$\begin{cases} M_{[ab]}^i = 0, & M_{a[bc]}^i = -R_{jpq}^i M_a^p M_b^q M_c^j, \\ M_{ab[cd]}^i = -R_{pqj}^i M_{ab}^p M_c^q M_d^j. \end{cases}$$

The values M_a^i , M_{ab}^i and M_{abc}^i form a fundamental third-order difference-geometric object of the surface \mathfrak{N} .

The normal vector of the hypersurface \mathfrak{N} at the point T satisfies the equations

$$g_{ij} n^i M_a^j = 0, \quad g_{ij} n^i n^j = 1.$$

A metric tensor of the hypersurface \mathfrak{N} is written in the form

$$g_{ab} = g_{ij} M_a^i M_b^j$$

and $\nabla g_{ab} = g_{abc} \psi^c$, where $g_{abc} = g_{ij} M_a^i M_{bc}^j + g_{ij} M_{ac}^i M_b^j$.

The vectors $M_{ab}^i e_i$ and $n_a^i e_i$ admit representations in the form of a linear combination of vectors of the reference point $\{T, M_a, n\}$:

$$M_{ab}^i e_i = Q_{ab}^c M_c + \mathcal{L}_{ab} n, \quad (4)$$

$$n_a^i e_i = \mathcal{L}_a^b M_b + n_a n, \quad (5)$$

where

$$Q_{ab}^c = g^{cd} g_{ik} M_{ab}^i M_d^k, \quad \mathcal{L}_{ab} = g_{ki} n^k M_{ab}^i, \quad \mathcal{L}_a^b = -g^{cb} \mathcal{L}_{ca}, \quad n_a = g_{ki} n^k n_a^i. \quad (6)$$

We call equations (4) and (5) the Gauss-Weingarten formulas of the hypersurface \mathfrak{N} . From (6) we obtain

$$\nabla \mathbb{Q}_{ab}^c + M_{ab}^c = \mathbb{Q}_{abd}^c \psi^d, \tag{7}$$

$$\nabla \mathcal{L}_{ab} = \mathcal{L}_{abc} \psi^c, \tag{8}$$

where

$$\begin{cases} \mathbb{Q}_{abd}^c = g^{ec} g_{ik} M_{ab}^i M_e^k + g^{ei} g_{ik} M_{abd}^i M_e^k + g^{ec} g_{ik} M_{ab}^i M_{ed}^k, \\ \mathcal{L}_{abc} = g_{ki} n_c^k M_{ab}^i + g_{ki} n^k M_{abc}^i, \\ g^{ec}_{\dots,d} = -g^{ea} g^{bc} g_{abd}. \end{cases}$$

From (7) and (8) it follows that \mathbb{Q}_{ab}^c is the object of affine connectedness and \mathcal{L}_{ab} is the tensor.

We call the object \mathbb{Q}_{ab}^c the object of induced affine connectedness of the hypersurface \mathfrak{N} . It is easy to prove that the induced affine connectedness and the internal affine connectedness coincide. The 1-forms of this connectedness have the form

$$\tilde{\psi}_b^a = \psi_b^a + \mathbb{Q}_{bc}^a \psi^c.$$

It is obvious that

$$D\psi^a = \psi^b \wedge \tilde{\psi}_b^a, \quad D\tilde{\psi}_b^a = \tilde{\psi}_b^c \wedge \tilde{\psi}_c^a + M_{bcd}^a \psi^c \wedge \psi^d,$$

where

$$M_{bcd}^a = \mathbb{Q}_{b[cd]}^a - \mathbb{Q}_{e[c}^a \mathbb{Q}_{b]d}^e.$$

The values M_{bcd}^a form the tensor which we call the curvature tensor of the hypersurface \mathfrak{N} . By extending equation (4) we obtain

$$R_{j\dot{p}q}^i M_a^k M_b^p M_c^q = (M_{abc}^d - \mathcal{L}_{a[b} \mathcal{L}_{c]}^d) M_d^i - (\nabla_{[c}^k \mathcal{L}_{|a|b]} - M_{bc}^d \mathcal{L}_{ad} + \mathcal{L}_{a[b} n_{c]}) n^i,$$

where ∇_c^k is the symbol of nonholomorphic covariant differentiation.

From the above equalities we obtain the generalized Gauss equations

$$R_{ipq} M_a^j M_b^p M_c^q M_e^i = M_{abce}^i + \mathcal{L}_{a[b} \mathcal{L}_{c]d} \tag{9}$$

and the generalized Peterson–Codazzi–Mainardi equations

$$R_{kpqr} M_a^k M_b^q M_c^r n^p = M_{bc}^d \mathcal{L}_{ad} - \nabla_{[c}^k \mathcal{L}_{|a|b]} - \mathcal{L}_{a[b} n_{c]}, \tag{10}$$

where

$$R_{ipqr} = q_{ij} R_{pqr}^j, \quad M_{abce} = g_{de} M_{abc}^d.$$

Equations (9) and (10) establish the connection between the curvature tensor of the space $Lm(Vn)$ and the curvature tensor of the hypersurface \mathfrak{N} in $Lm(Vn)$ [1, 5].

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