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# PERIODIC SOLUTIONS FOR FOURTH-ORDER DIFFERENTIAL EVOLUTION EQUATION INVOLVING POLY-HARMONIC OPERATOR

**Abstract.** In this note, we analyze the existence and uniqueness of periodic solutions for a fourthorder evolution differential equation involving the well known poly-harmonic operator. The right-hand term of the equation is taken in some anistropic Hölder spaces. Our approach is based on the study of a fourth- order abstract differential equation. To this end, we opt for the use of analytic semigroup's techniques.

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### 1 Motivation and statement of the problem

The theory of boundary value problems involving the poly-harmonic operator is a well-developed subject. For more information, we can refer the reader to [1-4, 17]. These works discuss a classical situation and establish the elliptic theory for the following model problem

$$\Delta^m u = f. \tag{1.1}$$

This differential equation is called the Kirchhoff-Love model for the vertical deflection of a thin elastic plate. The investigation of equation (1.1) was considered under different boundary conditions. The techniques of investigation are based on the use of the well-known potential theory or via the variational techniques. For the reader convenience, we just recall the classical poly-harmonic operator  $\Delta^m$  which can be regarded as iterations of the Laplace operator, that is,

$$\Delta^m = \Delta(\Delta^{m-1}).$$

On the other hand, there exists another research axis which are concerned with the study of evolution equations involving a poly-harmonic operator. For instance, we quote two famous model equations

$$D_t u + (-\Delta)^m u = f,$$

and

$$D_t^2 u + (-\Delta)^m u = f.$$

These classes of problems have been well investigated in different contexts and under several conditions. For more details, we refer the readers to [6, 20] and the references therein. All these studies were motivated by the fact that this kind of problems arise in several models describing various phenomena in the applied sciences. In our situation, we deal with the solvability of Cauchy problems for the fourth-order evolution equations involving a relaxed poly-harmonic operator. More precisely, we provide a systematic study of the following equation:

$$\left(D_t + (-1)^m L_m\right)^4 u(t, x) = f(t, x), \quad (t, x) \in \mathbb{R}^+ \times \Omega, \quad m \in \mathbb{N} - \{0\},$$
(1.2)

where  $L_m$  is the higher order differential operator defined by

$$L_m = \left(\sum_{i=1}^2 D_{x_i}^2 - \lambda\right)^m, \ \lambda > 0,$$

Here,  $x = (x_1, x_2)$  is a generic point of  $\mathbb{R}^2$  and  $\Omega = [0, 1] \times [0, 1]$ . The right-hand term of equation (1.2) belongs to the anisotropic Hölder space  $C^{\theta}(\mathbb{R}^+; L^p(\Omega))$  with  $0 < \theta < 1$ ,  $1 , consisting of all <math>\theta$ -Hölder continuous functions  $f : \mathbb{R}^+ \to L^p(\Omega)$  such that

$$\|f\|_{C^{\theta}(\mathbb{R}^{+};L^{p}(\Omega))} := \|f\|_{C_{b}(\mathbb{R}^{+};L^{p}(\Omega))} + \sup_{t \in \mathbb{R}^{+}} \frac{\|f(t) - f(\tau)\|_{L^{p}(\Omega)}}{|t - \tau|^{\theta}} < \infty,$$

where

$$C_b(\mathbb{R}^+; L^p(\Omega)) = \Big\{ f \in C(\mathbb{R}^+; L^p(\Omega)) : \lim_{t \to +\infty} f(t, x) < +\infty \Big\}.$$

We also assume that f is a T-periodic function in the first variable

$$f(t+T,\,\cdot\,) = f(t,\,\cdot\,), \ t \in \mathbb{R}^+.$$

Let us supplement Eq. (1.2) with the following higher order boundary conditions:

$$\begin{cases} [L_m]^k \Big[ \sum_{j=0}^1 (-1)^{j+1} D_{x_1}^j \Big]^m u \Big|_{\mathbb{R}^+ \times \{0\} \times [0,1]} = 0, \\ [L_m]^k u \Big|_{\mathbb{R}^+ \times \{1\} \times [0,1]} = 0, \end{cases}$$
(1.3)

and

$$[L_m]^k \left[\sum_{j=0}^2 D_{x_2}^j u\right]^m \bigg|_{\mathbb{R}^+ \times [0,1] \times \{1\}} = 0, \quad [L_m]^k \left[\sum_{j=0}^2 D_{x_2}^j u\right]^m \bigg|_{\mathbb{R}^+ \times [0,1] \times \{0\}} = 0, \tag{1.4}$$

where  $k \in \{0, 1, 2, 3\}$ . We also impose the following initial conditions of periodic type:

$$D_t^j u \big|_{\{0\} \times \Omega} = D_t^j u \big|_{\{T\} \times \Omega}, \quad j = 0, \dots, 3.$$
(1.5)

The boundary conditions (1.4) involve the well known Wentzell type boundary conditions which were introduced in [22] in order to study a multidimensional Diffusion Processes. Note that in this work, the diversity of boundary and initial conditions imposed on equation (1.2) and the character of the functional framework make the use of classical approaches a difficult task. For these reasons, we will opt for the use of an abstract point of view. As in [13], the abstract version of problem (1.2)-(1.5)will be treated by using the theory of analytic semigroups. This allows us to obtain some interesting regularity results for our problem.

## 2 Operational formulation of problem (1.2)-(1.5)

We set  $E = L^p(\Omega)$  endowed with its natural norm and define the vector-valued functions

$$u: \mathbb{R}^+ \to E; \ t \to u(t); \ u(t)(x) = u(t,x),$$
  
$$f: \mathbb{R}^+ \to E; \ t \to f(t); \ f(t)(x) = f(t,x).$$

Consider the operator  $\mathcal{A}_m$  defined by

$$\mathcal{A}_m \varphi(x) = ((-1)^m L_m) \varphi(x) \tag{2.1}$$

and its natural domain

$$D(\mathcal{A}_m) = \left\{ \varphi, \ [L_m]\varphi \in E : \ \varphi \Big|_{\{1\} \times [0,1]} = 0 \text{ and} \\ \left( \sum_{j=0}^1 (-1)^{j+1} D_{x_1}^j \right)^m \varphi \Big|_{\{0\} \times [0,1]} = 0, \ \left( \sum_{j=0}^2 D_{x_2}^j \right)^m \varphi \Big|_{[0,1] \times \{0,1\}} = 0 \right\}.$$

Then our problem (1.2) can be formulated as a complete fourth-order abstract differential equation

$$\begin{cases} u^{(4)}(t) + \sum_{k=1}^{3} {4 \choose k} \mathcal{A}_{m}^{k} u^{(4-k)}(t) + \mathcal{A}_{m}^{4} u(t) = f(t), & t \in \mathbb{R}^{+}, \\ u^{(j)}(0) = u^{(j)}(T), & 1 \le j \le 3, & T > 0, \end{cases}$$

$$(2.2)$$

with

$$\mathcal{A}_m^k = \mathcal{A}_m(\mathcal{A}_m^{k-1}).$$

To establish our main results, we need to use some spectral properties of the operator  $\mathcal{A}_m$  defined by (2.1). It is obvious that this study needs to introduce the operator  $(\mathcal{A}_0, D(\mathcal{A}_0))$  defined as follows:

$$\mathcal{A}_0\varphi(x) = \Big(\sum_{i=1}^2 D_{x_i}^2 - \lambda\Big)\varphi(x), \quad x = (x_1, x_2) \in \Omega, \tag{2.3}$$

and its natural domain

$$D(\mathcal{A}_0) = \left\{ \varphi, \sum_{i=1}^2 D_{x_i}^2 \varphi \in E : \\ \varphi \Big|_{\{0\} \times [0,1]} = 0, \sum_{j=0}^1 (-1)^{j+1} D_{x_1}^j \varphi \Big|_{\{1\} \times [0,1]} = 0, \sum_{j=0}^2 D_{x_2}^j \varphi \Big|_{[0,1] \times \{0,1\}} = 0 \right\}.$$

For  $f \in C^{\theta}(\mathbb{R}^+; L^p(\Omega))$ , we search for a strict solution u of problem (2.2), that is, a vectorial function u such that

$$\begin{cases} (i) & u \in C_b^4(\mathbb{R}^+; E) \cap C_b(\mathbb{R}^+; D(\mathcal{A}_m^4)), \\ (ii) & u^{(4-k)} \in C_b(\mathbb{R}^+; D(\mathcal{A}_m^k)), \quad k = 1, 2, 3 \end{cases}$$

## 3 On the elliptic character of the operator $\mathcal{A}_m$

The investigation of spectral properties of operator (2.3) is based on the study of the following spectral problem:

$$\begin{cases} D_{x_1}^2 v(x_1, x_2) + D_{x_2}^2 v(x_1, x_2) - \lambda v(x_1, x_2) = \varphi(x_1, x_2), & (x_1, x_2) \in \Omega, \\ v\big|_{x_1=1} = 0, & \sum_{j=0}^1 (-1)^{j+1} D_{x_1}^j v\big|_{x_1=0}, & \sum_{j=0}^2 D_{x_2}^j v(\cdot, 0) = \sum_{j=0}^2 D_{x_2}^j v(\cdot, 1) = 0. \end{cases}$$
(3.1)

The study of (3.1) is performed in the Lebesgue space  $L^p(\Omega)$ . As in [7], we use the commutative version of the well-known sum's operator theory developed by G. da Prato and P. Grisvard in [15]. Briefly, the abstract version of (3.1) can be formulated as follows:

$$\mathcal{M}v + \mathcal{N}v - \lambda v = \varphi, \ v \in D(\mathcal{M}) \cap D(\mathcal{N}),$$

where  $(\mathcal{M}, D(\mathcal{M}))$  and  $(\mathcal{N}, D(\mathcal{N}))$  are the linear operators defined as

$$D(\mathcal{M}) = \left\{ v, \ D_{x_1}^2 v \in L^p(0,1) : \ D_{x_1} v \big|_{x_1=0} - v \big|_{x_1=0} = 0, \ v \big|_{x_1=1} = 0 \right\}, \ \mathcal{M}v = D_{x_1}^2 v(x_1,x_2)$$
(3.2)

and

$$D(\mathcal{N}) = \left\{ v, \ D_{x_2}^2 v \in L^p(0,1) : \ \sum_{j=0}^2 D_{x_2}^j v \Big|_{x_2=0} = \sum_{j=0}^2 D_{x_2}^j v \Big|_{x_2=1} = 0 \right\}, \ \mathcal{N}v = D_{x_2}^2 v(x_1,x_2).$$

The basic spectral properties of the operator  $(\mathcal{M}, D(\mathcal{M}))$  are summarized in the following

**Lemma 3.1.** Let  $(\mathcal{M}, D(\mathcal{M}))$  be the linear operator defined by (3.2). Then  $\mathcal{M}$  is a closed linear operator satisfying the following properties:

(1)  $\mathbb{R}^+ \supset \rho(\mathcal{M}) \text{ and } \exists C > 0 : \forall \mu \ge 0,$ 

$$\|(\mathcal{M} - \mu I)^{-1}\|_{L(E)} \le \frac{C}{1+\mu},$$
(3.3)

 $(\rho(\mathcal{M}) \text{ is the resolvent set of } \mathcal{M});$ 

(2) 
$$\overline{D(\mathcal{M})} = L^p(0,1).$$

Proof. As in [19, p. 89], an explicit calculus shows that

$$\left\{ (\mathcal{M} - \mu I)^{-1} \varphi \right\}(x_1) = -\int_0^1 \mathcal{K}(\mu, s, x) \varphi(s) \, ds,$$

where  $\varphi \in L^p(0,1)$  and

$$\mathcal{K}(\mu, s, x) = \begin{cases} \frac{\sinh\sqrt{\mu}\,(1-x_1)[\sinh\sqrt{\mu}\,s + \sqrt{\mu}\,\cosh\sqrt{\mu}\,s]}{\sqrt{\mu}\,[\sinh\sqrt{\mu} + \sqrt{\mu}\,\cosh\sqrt{\mu}]}\,, & 0 \leqslant s \leqslant x_1, \\ \frac{\sinh\sqrt{\mu}\,(1-s)[\sinh\sqrt{\mu}\,x_1 + \sqrt{\mu}\,\cosh\sqrt{\mu}\,x_1]}{\sqrt{\mu}\,[\sinh\sqrt{\mu} + \sqrt{\mu}\,\cosh\sqrt{\mu}]}\,, & x_1 \leqslant s \leqslant 1. \end{cases}$$

Here,  $\sqrt{\mu}$  is the analytic determination defined by  $\operatorname{Re}\sqrt{\mu} > 0$ . Note that

$$\big|\sinh\sqrt{\mu} + \sqrt{\mu}\cosh\sqrt{\mu}\,\big| = \Big|\frac{e^{\operatorname{Re}\sqrt{\mu}}}{2}\,(a_{\sqrt{\mu}} + ib_{\sqrt{\mu}}) + \frac{e^{-\operatorname{Re}\sqrt{\mu}}}{2}\,(c_{\sqrt{\mu}} + id_{\sqrt{\mu}})\Big|,$$

where

$$\begin{cases} a_{\sqrt{\mu}} = 1 + \operatorname{Re} \sqrt{\mu} \, \cos \operatorname{Im} \sqrt{\mu} - \operatorname{Im} \sqrt{\mu} \, \sin \operatorname{Im} \sqrt{\mu} \,, \\ b_{\sqrt{\mu}} = 1 + \operatorname{Re} \sqrt{\mu} \, \sin \operatorname{Im} \sqrt{\mu} + \operatorname{Im} \sqrt{\mu} \, \cos \operatorname{Im} \sqrt{\mu} \,, \\ c_{\sqrt{\mu}} = (\operatorname{Re} \sqrt{\mu} - 1) \cos \operatorname{Im} \sqrt{\mu} - \operatorname{Im} \sqrt{\mu} \, \sin \operatorname{Im} \sqrt{\mu} \,, \\ d_{\sqrt{\mu}} = (1 - \operatorname{Re} \sqrt{\mu}) \sin \operatorname{Im} \sqrt{\mu} + \operatorname{Im} \sqrt{\mu} \, \cos \operatorname{Im} \sqrt{\mu} \,. \end{cases}$$

So, we deduce that

$$\begin{aligned} \left| \sinh \sqrt{\mu} + \sqrt{\mu} \cosh \sqrt{\mu} \right| \\ &\geq \frac{e^{\operatorname{Re}\sqrt{\mu}}}{2} \left[ \left( 1 + \operatorname{Re}\sqrt{\mu} \right)^2 + \left( \operatorname{Im}\sqrt{\mu} \right)^2 \right]^{1/2} - \frac{e^{-\operatorname{Re}\sqrt{\mu}}}{2} \left[ \left( 1 - \operatorname{Re}\sqrt{\mu} \right)^2 + \left( \operatorname{Im}\sqrt{\mu} \right)^2 \right]^{1/2} \\ &\geq \sinh \operatorname{Re}\sqrt{\mu} \left[ 1 + \left( \operatorname{Re}\sqrt{\mu} \right)^2 + 2\operatorname{Re}\sqrt{\mu} \right]^{1/2}, \end{aligned}$$

which implies

 $|\sinh\sqrt{\mu} + \sqrt{\mu}\cosh\sqrt{\mu}| \ge \sinh\operatorname{Re}\sqrt{\mu}\left[1 + \left(\operatorname{Re}\sqrt{\mu}\right)\right].$ 

To obtain the desired result, it suffices to observe that

$$\begin{aligned} \frac{\sinh(1-x_1)\operatorname{Re}\sqrt{\mu}}{\sinh\operatorname{Re}\sqrt{\mu}\left[1+(\operatorname{Re}\sqrt{\mu})\right]} &= \frac{e^{\operatorname{Re}\sqrt{\mu}(1-x_1)}-e^{-\operatorname{Re}\sqrt{\mu}(1-x_1)}}{(e^{\operatorname{Re}\sqrt{\mu}}-e^{-\operatorname{Re}\sqrt{\mu}})(1+(\operatorname{Re}\sqrt{\mu}))} \\ &= \frac{e^{-\operatorname{Re}\sqrt{\mu}x_1}[1-e^{-\operatorname{Re}\sqrt{\mu}(2-x_1)}]}{(1-e^{-\operatorname{Re}\sqrt{\mu}})(1+(\operatorname{Re}\sqrt{\mu}))} \leq \frac{Ce^{-\operatorname{Re}\sqrt{\mu}x_1}}{(1+(\operatorname{Re}\sqrt{\mu}))}, \end{aligned}$$

from which it follows that

$$\int_{0}^{1} \left| \frac{\sinh(1-x_{1})\operatorname{Re}\sqrt{\mu}}{\sinh\operatorname{Re}\sqrt{\mu}\left[1+(\operatorname{Re}\sqrt{\mu})\right]} \right|^{p} dx_{1} \leq \int_{0}^{1} |e^{-\operatorname{Re}\sqrt{\mu}x_{1}}|^{p} dx_{1} \leq \frac{1}{p\operatorname{Re}\sqrt{\mu}}.$$

The density of  $D(\mathcal{M})$  is obtained thanks to Proposition 1.1 in [18, p. 18].

Concerning the operator  $(\mathcal{N}, D(\mathcal{N}))$ , one has

**Lemma 3.2.** Let  $(\mathcal{N}, D(\mathcal{N}))$  be the linear operator defined by (3.2). Then  $\mathcal{N}$  is a closed linear operator satisfying the following conditions:

(1) 
$$\mathbb{R}^+ \supset \rho(\mathcal{N}) \text{ and } \exists C > 0 : \forall \mu \ge 0,$$
  
 $\| (\mathcal{N} - \mu I)^{-1} \|_{L(E)} \le \frac{C}{1 + \mu};$  (3.4)

(2)  $\overline{D(\mathcal{N})} = L^p(0,1).$ 

*Proof.* These results use the same argument as in the previous lemma. In fact, using the classical variation of a constant, we show that

$$\left\{ (\mathcal{N} - \mu I)^{-1} \varphi \right\}(x_2) = \frac{1}{(1 - e^{-2\sqrt{\mu}})} \int_0^1 e^{-\sqrt{\mu}(s - x_2 - 2)} \varphi(s) \, ds + \frac{1}{(1 - e^{-2\sqrt{\mu}})} \int_0^1 e^{-\sqrt{\mu}(2 - s + x_2)} \varphi(s) \, ds$$

$$+ \frac{\mathcal{C}(\mu)}{(1 - e^{-2\sqrt{\mu}})} \int_{0}^{1} e^{-\sqrt{\mu}(2 - s - x_2)} \varphi(s) \, ds + \frac{\mathcal{S}(\mu)}{(1 - e^{-2\sqrt{\mu}})} \int_{0}^{1} e^{-\sqrt{\mu}(s + x_2)} \varphi(s) \, ds \\ + \int_{0}^{1} e^{-\sqrt{\mu}(s - x_2)} \varphi(s) \, ds + \frac{1}{2\sqrt{\mu}} \int_{0}^{x_2} e^{-\sqrt{\mu}(x_2 - s)} \varphi(s) \, ds + \frac{1}{2\sqrt{\mu}} \int_{0}^{x_2} e^{-\sqrt{\mu}(s - x_2)} \varphi(s) \, ds,$$

where

$$\mathcal{C}(\mu) = \frac{(\mu - \sqrt{\mu} + 1)}{(\mu + \sqrt{\mu} + 1)}$$
 and  $\mathcal{S}(\mu) = \frac{1}{\mathcal{C}(\mu)}$ .

At this level, a direct computation shows that

$$\|(\mathcal{N} - \mu I)^{-1}\|_{L(E)} \le \frac{C}{1 + |\mu|}.$$

Now, we have

**Proposition 3.1.** Let  $(\mathcal{A}_0, D(\mathcal{A}_0))$  be the linear operator defined by (2.3). Then  $\mathcal{A}_0$  is a closed linear densely defined operator satisfying the natural ellipticity hypothesis

$$\mathbb{R}^+ \subset \rho(\mathcal{A}_0) \quad and \quad \exists C > 0, \quad \forall z \ge 0 \quad \|(\mathcal{A}_0 - zI)^{-1}\| \le \frac{C}{1+z} \,. \tag{3.5}$$

*Proof.* As a direct consequence of the use of the commutative version of the sum operators technique due to [15], we conclude that the operator  $(\mathcal{A}_0 - zI)^{-1}$  is well defined and

$$(\mathcal{A}_0 - zI)^{-1}\varphi = -\frac{1}{2i\pi} \int_{\Gamma} (\mathcal{N} + zI)^{-1} (\mathcal{M} - zI - \lambda I)^{-1} \varphi \, dz, \qquad (3.6)$$

where  $\Gamma$  is a suitable Jordan curve lying in  $\sigma(-\mathcal{N}) \cap \sigma(\mathcal{M})$ . Keeping in mind estimates (3.3) and (3.4), estimate (3.5) can be easily deduced from formula (3.6). Furthermore, this estimate holds true if we replace z by  $z + \lambda$ .

As a consequence of these results, we have

**Lemma 3.3.** The operator  $(\mathcal{A}_m, D(\mathcal{A}_m))$  defined by (2.1) is closed, densely defined and satisfies the Krein-ellipticity property, that is,  $\mathbb{R}^+ \subseteq \rho(\mathcal{A}_m)$  and there exists C > 0 such that for all  $z \ge 0$  we have

$$\|(\mathcal{A}_m - zI)^{-1}\|_{\mathcal{L}(L^p(\Omega))} \le \frac{C}{1+z}.$$
 (3.7)

*Proof.* The proof is based on the use of the following algebraic formula:

$$(\mathcal{A}_m - \lambda I)^{-1} = \lambda^{1 - \frac{1}{m}} \sum_{k=0}^{m-1} b_k \left( (\mathcal{A}_0 - \varepsilon_k \lambda^{\frac{1}{m}} I)^{-1} \right)$$

with

$$b_k = \frac{1}{\prod_{n=0}^{m-1} (\varepsilon_k - \varepsilon_n)}, \quad n \neq k,$$

where  $\varepsilon_k$  takes all the values of the *m*th root of unity. By adapting the same techniques used in [21], we conclude that for  $\lambda > 0$ ,

$$\left\| (\mathcal{A}_m - \lambda I)^{-1} \right\|_E = \mathcal{O}\left(\frac{1}{\lambda}\right).$$

**Remark 3.1.** Using the classical argument of continuation of the resolvent, we know that estimate (3.7) holds true in a sector of the form

$$\Sigma = \left\{ z \in \mathbb{C}^* : |\arg z| \le \pi \left( 1 - \frac{1}{2^m} \right) \right\} \cup B(0, \epsilon_0)$$

with some small  $\epsilon_0 > 0$ .

# 4 Construction of the solution of (2.2)

Differently from [9] and [10], we use another approach based essentially on the semigroup theory. In fact, from [5] we know that estimate (3.7) implies that the operator

$$\mathcal{B} = -(-\mathcal{A}_m)^{1/4}$$

is well defined and it is the infinitesimal generator of a generalized analytic semigroups  $(e^{t\mathcal{B}})_{\xi>0}$ . More precisely, there exists a sector

$$\Pi_{\delta, r_0} = \left\{ z \in \mathbb{C}^* : |\arg z| \le \delta + \frac{\pi}{2} \right\} \cup \overline{B(0, r_0)}$$

(with some positive  $\delta$ ,  $r_0$ ) and C > 0 such that

$$\rho(\mathcal{B}) \supset \Pi_{\delta, r_0} \text{ and } \forall z \in \Pi_{\delta, r_0}, \|(\mathcal{B} - zI)^{-1}\| \le \frac{C}{1 + |z|}.$$

Thus, for all t > 0 and  $\varphi \in E$ , one has

$$e^{t\mathcal{B}}\varphi = \frac{1}{2i\pi}\int\limits_{\gamma}e^{zt}(\mathcal{B}-zI)^{-1}\varphi\,dz,$$

where  $\gamma = \partial \Pi_{\delta, r_0}$  (the sectorial boundary curve of  $\Pi_{\delta, r_0}$  oriented from  $\infty e^{i(\delta + \pi/2)}$  to  $\infty e^{-i(\delta + \pi/2)}$ ). Then we have the following auxiliary results:

(1)  $\forall k \in \mathbb{N}, \exists m_k \ge 1, \omega > 0$ :

$$\|t^k \mathcal{B}^k e^{t\mathcal{B}}\|_{L(E)} \le m_k e^{-\omega t},\tag{4.1}$$

(2)  $\lim_{t\to 0} e^{t\mathcal{B}}\varphi = \varphi$  if and only if  $\varphi \in \overline{D(\mathcal{B})}$ .

Using the same reasoning as in [13], the representation formula of the solution for problem (2.2) can be deduced from the scalar case. Putting

$$\Psi(\mathcal{B}) = (1 - e^{T\mathcal{B}})^{-1}$$

and

$$\Phi(\mathcal{B}) = -\frac{1}{2} \left( 1 + 4e^{\mathcal{B}T} + e^{2\mathcal{B}T} \right) T^3,$$

the formal solution of (2.2) is formulated as follows:

$$u(t) = \sum_{k=1}^{4} u_k(t) \tag{4.2}$$

with

$$\begin{split} u_1(t) &= \frac{\Psi(\mathcal{B})}{6} \int_t^{t+T} (s-t)^3 e^{(t+T-s)\mathcal{B}} f(s) \, ds, \\ u_2(t) &= -\frac{\Psi^2(\mathcal{B})}{2} T \int_t^{t+T} (s-t)^2 e^{(t+T-s)\mathcal{B}} f(s) \, ds, \\ u_3(t) &= \frac{\Psi^3(\mathcal{B})}{2} T^2 \int_t^{t+T} (s-t) e^{(t+T-s)\mathcal{B}} f(s) \, ds, \\ u_4(t) &= -\frac{\Phi(\mathcal{B})\Psi^4(\mathcal{B})}{2} T^3 \int_t^{t+T} e^{(t+T-s)\mathcal{B}} f(s) \, ds. \end{split}$$

First of all, we have

**Proposition 4.1.** Let  $f \in C^{\theta}(\mathbb{R}^+; E)$  with  $\theta \in [0, 1[$ . Then, for all  $t \ge 0$  and  $k \in \mathbb{N}$ ,

$$u(t) \in D(\mathcal{B}^k).$$

*Proof.* It suffices to show that

$$\|\mathcal{B}^k u(t)\|_{C^{\theta}(\mathbb{R}^+;E)} < \infty,$$

which means that for i = 1, 2, 3, 4, we have

$$\|\mathcal{B}^k u_i(t)\|_{C^{\theta}(\mathbb{R}^+;E)} < \infty.$$

Since these vectorial functions can be treated similarly, we restrict ourselves on the first one. First, recall that

$$u_1(t) = \frac{e^{(t+T)\mathcal{B}}\Psi(\mathcal{B})}{6} \int_{t}^{t+T} (s-t)^3 e^{-\mathcal{B}s} f(s) \, ds,$$

since

$$e^{-t\mathcal{B}} \subset \bigcap_{k\geq 1} D(\mathcal{B}^k),$$

and the Banach valued functions  $\Psi$  and  $\Phi$  are regular in the sense that there are bounded quantities. Then it follows that

$$\mathcal{B}^{k}u_{1}(t) = \frac{\mathcal{B}^{k}e^{(t+T)\mathcal{B}}}{6}\Psi(\mathcal{B})\bigg(\int_{t}^{t+T}(s-t)^{3}e^{-\mathcal{B}s}(f(s)-f(t))\,ds - \int_{t}^{t+T}(s-t)^{3}\mathcal{B}^{k}e^{-\mathcal{B}s}f(t)\,ds\bigg),$$

and, clearly,

$$\begin{aligned} \|\mathcal{B}^{k}u_{1}(t)\|_{E} &\leq C \bigg( \int_{0}^{+\infty} (s-t)^{3+\theta} e^{-\omega s} \, ds \bigg) \|f\|_{C^{\theta}(\mathbb{R}^{+};E)} + C' \|f(t)\|_{E} \\ &\leq C\Gamma(2+\theta) \|f\|_{C^{\theta}(\mathbb{R}^{+};E)} + C' \|f(t)\|_{E} \leq C \|f\|_{C^{\theta}(\mathbb{R}^{+};E)}, \end{aligned}$$

where  $\Gamma$  is the usual Euler function defined by

$$\Gamma(z) = \int_{0}^{+\infty} e^{-w} w^{z-1} dw, \quad \operatorname{Re} z > 0.$$

The study of the regularity of solution (4.2) is purely technical. As for the techniques, we refer the reader to [8, 11-14, 16, 16] and the references therein, in which the kernel semigroups appearing in formula (4.2) were extensively used in different situation to provide some interesting results for different abstract differential equation under abstract boundary conditions.

**Proposition 4.2.** Let  $f \in C^{\theta}(\mathbb{R}^+; E)$  with  $\theta \in ]0, 1[$ . Then:

- (1) for all  $j \in \{1, 2, 3\}$ :  $u^{(j)} \in C_b(\mathbb{R}^+; E)$ .
- (2) for all  $k \in \{0, 1, 2, 3\}$ :  $\mathcal{A}_m^{4-k} u \in C_b(\mathbb{R}^+; E)$ .

Summing up, we are able to state our main regularity results for Problem (2.2).

**Proposition 4.3.** Let  $f \in C^{\theta}(\mathbb{R}^+; E)$  with  $\theta \in ]0,1[$ . Then Problem (2.2) has a unique solution

$$\begin{cases} (i) & u \in C_b^4(\mathbb{R}^+; E) \cap C_b(\mathbb{R}^+; D(A^4)), \\ (ii) & u^{(j)} \in C_b(\mathbb{R}^+; D(A^{4-j})), \quad j \in \{1, 2, 3\} \end{cases}$$

Moreover, the solution u satisfies the following maximal regularity property

$$u^{(4)}, \mathcal{A}_m^j u^{(4-j)}, \mathcal{A}_m^4 u \in C^{\theta}(\mathbb{R}^+; E).$$

Now, we are able to give our main regularity results for our concrete problem (1.2)–(1.5) which can be formulated as follows.

**Proposition 4.4.** Let  $f \in C^{\theta}(\mathbb{R}^+; L^p(\Omega))$  with  $0 < \theta < 1$ , 1 . Then Problem (1.2)–(1.5) has a unique strict solution

 $u \in C_b^4(\mathbb{R}^+; L^p(\Omega))$ 

such that for all  $(t, x) \in \mathbb{R}^+ \times \Omega$ , we have

- (1)  $D_{x_1}^{2m}u(t,x), D_{x_2}^{2m}u(t,x) \in C_b(\mathbb{R}^+; W^{8m,p}(\Omega));$
- (2)  $D_t^4 u(t,x) \in C_b^4(\mathbb{R}^+; L^p(\Omega));$
- (3)  $D_t^{4-k}[D_{x_1}^2 + D_{x_2}^2]^k u(t,x) \in C_b(\mathbb{R}^+; L^p(\Omega)), \ k \in \{1,2,3\}.$

*Proof.* First, the use of the sum operator technique allows us to conclude that

$$D_{x_1}^2 u(\cdot, x), D_{x_2}^2 u(\cdot, x) \in L^p(\Omega),$$

which justify the first assertion. The assertions (2) and (3) are viewed as direct consequence of the preceding theorem.  $\Box$ 

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