

Boundary-Value Problems for Weakly Singular Integral Equations

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In the space $L_2[a, b]$, we consider a perturbed linear boundary-value problem for a weakly singular integral equation

$$x(t) - \int_a^b K(t, s)x(s) ds = f(t) + \varepsilon \int_a^b \overline{K}(t, s)x(s) ds, \quad (1)$$

$$lx(\cdot) = \alpha + \varepsilon Jx(\cdot). \quad (2)$$

We establish conditions for the bifurcation of solutions of the boundary-value problem (1), (2) and determine the structure of these solutions under the condition that the generating boundary-value problem

$$x(t) - \int_a^b K(t, s)x(s) ds = f(t), \quad lx(\cdot) = \alpha \quad (3)$$

is unsolvable.

Here, $K(t, s) = \frac{H(t, s)}{|t-s|^\gamma}$ and $\overline{K}(t, s) = \frac{\overline{H}(t, s)}{|t-s|^\beta}$, where $H(t, s), \overline{H}(t, s)$ are functions bounded in the domain $[a, b] \times [a, b]$, $0 < \gamma < 1$, $0 < \beta < 1$, $f \in L_2[a, b]$, $l = \text{col}(l_1, l_2, \dots, l_p) : L_2[a, b] \rightarrow \mathbb{R}^p$, $J = \text{col}(J_1, J_2, \dots, J_p) : L_2[a, b] \rightarrow \mathbb{R}^p$ are bounded linear functionals, $l_\nu, J_\nu : L_2[a, b] \rightarrow \mathbb{R}$, $\nu = \overline{1, p}$, $\alpha = \text{col}(\alpha_1, \alpha_2, \dots, \alpha_p) \in \mathbb{R}^p$ and $\varepsilon \ll 1$ is a small parameter.

By using the results obtained in [2], we show that the study of the problem of appearance of solutions of the boundary-value problem (1), (2) reduces to the corresponding task for the perturbed boundary-value problem for the Fredholm integral equation

$$x(t) = f_n(t) + \sum_{k=0}^n \varepsilon^k \int_a^b R_n^k(t, s)x(s) ds, \quad (4)$$

$$f_n(t) = f(t) + \sum_{k=1}^{n-1} \int_a^b R_k^0(t, s)f(s) ds + \sum_{k=1}^{n-1} \varepsilon^k \sum_{m=k}^{n-1} \int_a^b R_m^k(t, s)f(s) ds,$$

where $R_n^k(t, s)$, $k = \overline{0, n}$, are the sums of C_n^k kernels of all possible products of $n - k$ integral operators K and k integral operators \overline{K}

$$(Kw)(t) = \int_a^b \frac{H(t, s)}{|t-s|^\gamma} w(s) ds \quad \text{and} \quad (\overline{K}w)(t) = \int_a^b \frac{\overline{H}(t, s)}{|t-s|^\beta} w(s) ds.$$

We apply the approach described in [3] to the study of the boundary-value problem (4), (2) and show that it can be reduced to the operator equation. Let $\{\varphi_i(t)\}_{i=1}^\infty$ be a complete orthonormal

system of functions in $L_2[a, b]$. We introduce the notation

$$\begin{aligned}
 x_i &= \int_a^b x(t)\varphi_i(t) dt, \quad a_{ij} = \int_a^b \int_a^b K_n(t, s)\varphi_i(t)\varphi_j(s) dt ds, \\
 a_{ij}^k &= \int_a^b \int_a^b R_n^k(t, s)\varphi_i(t)\varphi_j(s) dt ds, \quad k = \overline{1, n}, \\
 f_i &= \int_a^b f(t)\varphi_i(t) dt + \sum_{k=1}^{n-1} \int_a^b \int_a^b R_k^0(t, s)f(s)\varphi_i(t) dt ds, \\
 f_i^k &= \sum_{m=k}^{n-1} \int_a^b \int_a^b R_m^k(t, s)f(s)\varphi_i(t) ds dt, \quad k = \overline{1, n-1}.
 \end{aligned}$$

By using this notation in the boundary-value problem (4), (2), we obtain the operator equation:

$$Uz = q + \sum_{k=1}^{n-1} \varepsilon^k q_k + \sum_{k=1}^n \varepsilon^k U_k z, \tag{5}$$

where

$$\begin{aligned}
 U &= \begin{bmatrix} \Lambda \\ W \end{bmatrix}, \quad U_1 = \begin{bmatrix} \Lambda_1 \\ W_1 \end{bmatrix}, \quad U_k = \begin{bmatrix} \Lambda_k \\ 0 \end{bmatrix}, \quad k = \overline{2, n}, \\
 q &= \begin{bmatrix} g \\ \alpha \end{bmatrix}, \quad q_k = \begin{bmatrix} g_k \\ 0 \end{bmatrix}, \quad k = \overline{1, n-1},
 \end{aligned}$$

where the vectors $z, g, g_k, k = \overline{1, n-1}$ and the matrices $W, W_1, \Lambda, \Lambda_k, k = \overline{1, n}$ have the form

$$\begin{aligned}
 z &= \text{col}(x_1, x_2, \dots, x_i, \dots), \quad g = \text{col}(f_1, f_2, \dots, f_i, \dots), \\
 g_k &= \text{col}(f_1^k, f_2^k, \dots, f_i^k, \dots), \quad W = l\Phi(\cdot), \quad W_1 = J\Phi(\cdot), \\
 \Lambda &= \begin{pmatrix} 1 - a_{11} & -a_{12} & \dots & -a_{1i} & \dots \\ -a_{21} & 1 - a_{22} & \dots & -a_{2i} & \dots \\ \dots & \dots & \dots & \dots & \dots \\ -a_{i1} & -a_{i2} & \dots & 1 - a_{ii} & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}, \quad \Lambda_k = \begin{pmatrix} a_{11}^k & a_{12}^k & \dots & a_{1i}^k & \dots \\ a_{21}^k & a_{22}^k & \dots & a_{2i}^k & \dots \\ \dots & \dots & \dots & \dots & \dots \\ a_{i1}^k & a_{i2}^k & \dots & a_{ii}^k & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}, \\
 \Phi(t) &= (\varphi_1(t), \varphi_2(t), \dots, \varphi_i(t), \dots).
 \end{aligned}$$

The generating equation for the operator equation (5) has the form

$$Uz = q. \tag{6}$$

The operator $\Lambda : \ell_2 \rightarrow \ell_2$ appearing on the left-hand side of the operator equation (6) has the form $\Lambda = I - A$, where $I : \ell_2 \rightarrow \ell_2$ is the identity operator and $A : \ell_2 \rightarrow \ell_2$ is a compact operator. Hence, according to S. Krein’s classification, the operator $\Lambda : \ell_2 \rightarrow \ell_2$ is a Fredholm operator of index zero ($\dim \ker \Lambda = \dim \ker \Lambda^* < \infty$) and the operator $U : \ell_2 \rightarrow \ell_2 \times \mathbb{R}^p$ is a Fredholm operator of nonzero index ($\dim \ker U < \infty, \dim \ker U^* < \infty$).

The following statement is true for equation (6) (see [4]).

Theorem 1. *The homogeneous equation (6) ($q = 0$) possesses a d_2 -parameter family of solutions $z \in \ell_2$,*

$$z = P_{\Lambda_r} P_{Q_{d_2}} c_{d_2} \quad \forall c_{d_2} \in \mathbb{R}^{d_2}.$$

The inhomogeneous equation (6) is solvable if and only if the following $r+d_1$ linearly independent conditions are satisfied:

$$P_{\Lambda_r^*} g = 0, \quad P_{Q_{d_1}^*} (\alpha - W\Lambda^+ g) = 0,$$

and the equation possesses a d_2 -parameter family of solutions $z \in \ell_2$ of the form

$$z = P_{\Lambda_r} P_{Q_{d_2}} c_{d_2} + P_{\Lambda_r} Q^+ (\alpha - W\Lambda^+ g) + \Lambda^+ g \quad \forall c_{d_2} \in \mathbb{R}^{d_2}.$$

Here, $Q = WP_{\Lambda_r}$ is a $(p \times r)$ -matrix, $P_{\Lambda_r} (P_{\Lambda_r^*})$ is a matrix formed by a complete system of r linearly independent columns (rows) of the matrix projector $P_{\Lambda} (P_{\Lambda^*})$, where $P_{\Lambda} (P_{\Lambda^*})$ is the projector onto the kernel (cokernel) of the matrix Λ , and $P_{Q_{d_2}} (P_{Q_{d_1}^*})$ is a matrix formed by the complete system of d_2 (d_1) linearly independent columns (rows) of the matrix projector $P_Q (P_{Q^*})$, where $P_Q (P_{Q^*})$ is the projector onto the kernel (cokernel) of the matrix Q and $\Lambda^+ (Q^+)$ is the pseudoinverse Moore–Penrose matrix for the matrix $\Lambda (Q)$.

We now determine the conditions required for the bifurcation of solutions of the perturbed inhomogeneous boundary-value problem (1), (2) and study the structure of these solutions under the conditions that the solution of the homogeneous generating the boundary-value problem (3) ($f(t) = 0, \alpha = 0$) is not unique, i.e. (see [2]), $P_{\Lambda_r} P_{Q_{d_2}} \neq 0$, and that the inhomogeneous generating boundary-value problem (3) is unsolvable.

It is known (see [9]) that small perturbations preserve the Fredholm property of the operator, i.e., the operator $(U - \sum_{k=1}^n \varepsilon^k U_k)$ is a Fredholm operator with nonzero index. This enables one to investigate equation (5) by the methods of the theory of perturbed operator boundary-value problems with Fredholm linear part (see, e.g., [1, 4, 11]) obtained as a generalization of the classical methods of the perturbation theory of periodic boundary-value problems in the theory of oscillations (see [5, 7, 8, 10]).

The analysis of the appearance of solutions of equation (5) is closely connected with the $((r + d_1) \times d_2)$ -matrix

$$B_0 = \begin{bmatrix} P_{\Lambda_r^*} \Lambda_1 P_{\Lambda_r} P_{Q_{d_2}} \\ P_{Q_{d_1}^*} (W_1 - W\Lambda^+ \Lambda_1) P_{\Lambda_r} P_{Q_{d_2}} \end{bmatrix},$$

constructed by using the coefficients of equation (5).

We introduce an $((r + d_1) \times (r + d_1))$ -matrix $P_{B_0^*}$, which is a projector onto the cokernel of the matrix B_0 and a matrix

$$G = \begin{bmatrix} -P_{\Lambda_r^*} & 0 \\ P_{Q_{d_1}^*} W\Lambda^+ & -P_{Q_{d_1}^*} \end{bmatrix},$$

formed by $r + d_1$ rows and infinitely many columns. Moreover, as the matrix B_0 , it is completely determined by the coefficients of equation (5).

By the Vishik–Lyusternik method (see [12]), we find efficient conditions for the coefficients guaranteeing the appearance of a family of solutions of the perturbed linear boundary-value problem (5) in the form of a Laurent series in powers of the small parameter ε with singularity at the point $\varepsilon = 0$.

The results obtained for the perturbed equations (5) enable us to establish the conditions for the existence of a d_2 -parameter family of solutions of the original perturbed boundary-value problem (1), (2). Indeed, if the boundary-value problem (1), (2) possesses at least one solution,

then, according to the Riesz–Fischer theorem, one can find an element $x \in L_2[a, b]$ such that the quantities x_i , $i = \overline{1, \infty}$, determined from equation (5) are the Fourier coefficients of this elements, i.e., the following representation is true:

$$x(t) = \Phi(t)z. \quad (7)$$

As in [6], we conclude that the element $x(t)$ given by relations (7) is the required d_2 -parameter family of solutions of the original boundary-value problem (1), (2). Therefore, the following statement is true.

Theorem 2. *Suppose that the generating boundary-value problem (3) is unsolvable. If conditions*

$$P_{\Lambda_r} P_{Q_{d_2}} \neq 0, \quad P_{B_0^*} G = 0,$$

are satisfied, then the boundary-value problem (1), (2) has a d_2 -parameter family of solutions in the form of series with singularity at the point $\varepsilon = 0$ convergent for sufficiently small fixed $\varepsilon \in (0, \varepsilon_)$.*

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